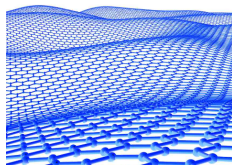
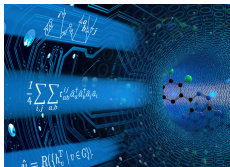


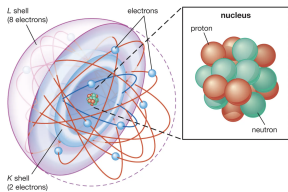
Mathematical foundations of modern quantum simulations



Fabian M. Faulstich

Department of Mathematics, Rensselaer Polytechnic Institute, Troy, NY

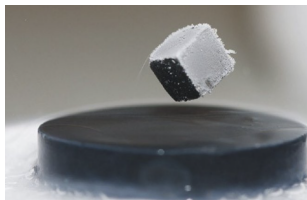
What are quantum simulations?



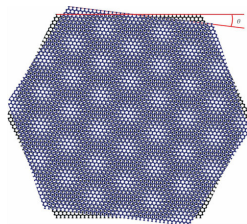
Stability of matter



Color of gold



Superconductivity

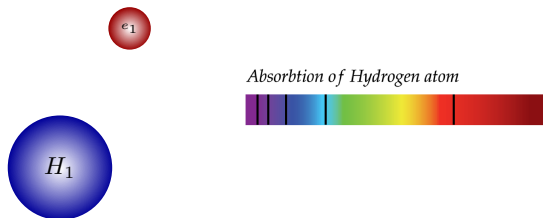


Solid state physics

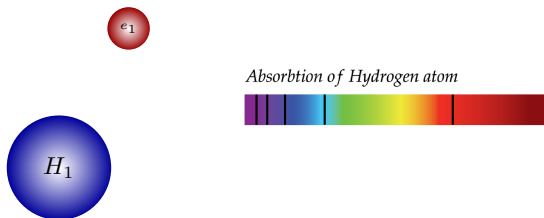
Towards a mathematical formulation



Towards a mathematical formulation



Towards a mathematical formulation



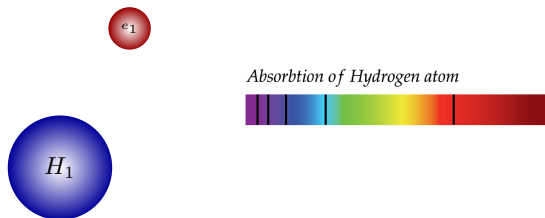
Two energetic contributions:

1. Kinetic energy: $h_{\text{kin}} = -\frac{1}{2}\Delta$

2. Coulomb energy: $V_{\text{coul}} = -\frac{1}{|r - R|}$



Towards a mathematical formulation



Two energetic contributions:

1. Kinetic energy: $h_{\text{kin}} = -\frac{1}{2}\Delta$

2. Coulomb energy: $V_{\text{coul}} = -\frac{1}{|r - R|}$



Birth of quantum chemistry^a:

$$\Delta\psi + \frac{2m}{\hbar^2} \left(E + \frac{e^2}{r} \right) \psi = 0$$

^aSchrödinger, Annalen der Physik (1926)

Schrödinger equation

The goal is to solve

$$H\Psi = E\Psi,$$

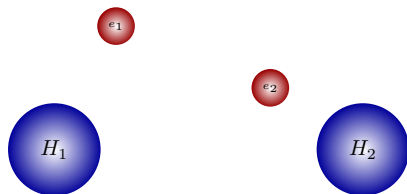
where

$$H = -\frac{1}{2} \sum_{i=1}^N \Delta_{r_i} - \sum_{i=1}^N \sum_{j=1}^{N_{\text{nuc}}} \frac{Z_j}{|r_i - R_j|} + \sum_{i=1}^N \sum_{j>i}^N \frac{1}{|r_i - r_j|}$$

and Ψ is a function defined on \mathbb{R}^{3N} for N electrons.

r_i position of the i th electron

R_j, Z_j position (fixed) and charge of the j th nucleus



What is the problem?

Discretization scales *exponentially* in the number of electrons N

¹The estimated number of particles in the universe is $\sim 10^{80}$

What is the problem?

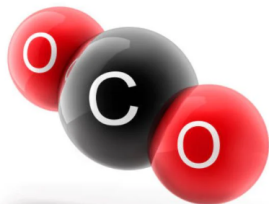
Discretization scales *exponentially* in the number of electrons N

Example: carbon dioxide (CO_2)

Oxygen has 8 electrons

Carbon has 6 electrons

In total: $N = 2 * 8 + 6 = 22$



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What is the problem?

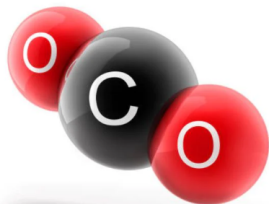
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Naïve grid with 10 points per axis yields $10^{3 \cdot 22} = 10^{66}$ grid points¹

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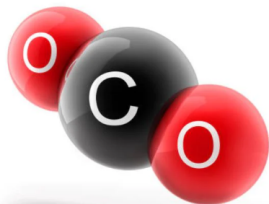
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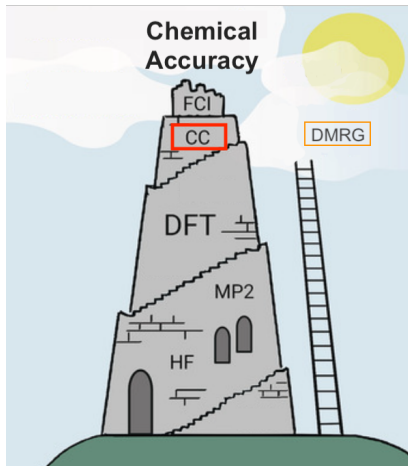


Naïve grid with 10 points per axis yields $10^{3 \cdot 22} = 10^{66}$ grid points¹

⇒ Approximation is key!

¹The estimated number of particles in the universe is $\sim 10^{80}$

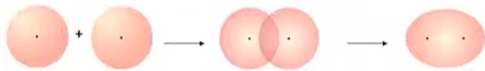
Quantum chemical methods



Strongly correlated quantum chemistry



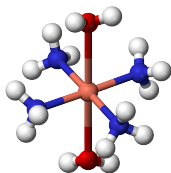
Strongly correlated quantum chemistry



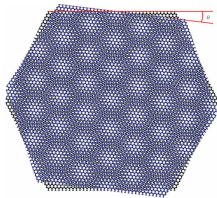
This happens at various scales



Small molecules

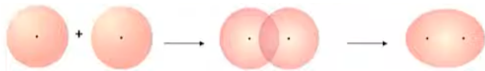


Transition metal
complexes



Twisted bilayer
graphene

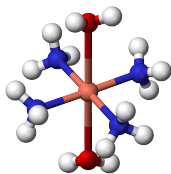
Strongly correlated quantum chemistry



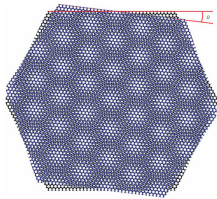
This happens at various scales



Small molecules



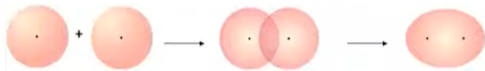
Transition metal
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Strongly correlated systems are *very* difficult to compute

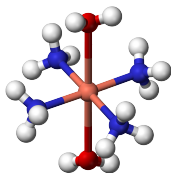
Strongly correlated quantum chemistry



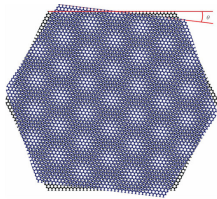
This happens at various scales



Small molecules



Transition metal
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Twisted bilayer
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Strongly correlated systems are *very* difficult to compute
Important failure modes for existing methods

⇒ The research frontier of computational chemistry lies in strongly correlated systems

Roadmap



1. Quantum mechanics of the hydrogen atom
2. The hydrogen molecule and the Hartree-Fock method
3. Second quantization I
4. Second quantization II
5. Second quantization III
6. The linear combination of atomic orbitals ansatz
7. Gaussian-type atomic orbitals and high-dimensional integration
8. Size consistency and the coupled cluster ansatz
9. Truncating the coupled cluster ansatz and its working equations
10. Finding one root to the coupled cluster equations
11. Finding all roots to the coupled cluster equations

Quantum mechanics of the hydrogen atom

Analytic solution to the eigenvalue problem

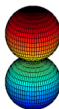
$$\left(\frac{1}{2} \Delta_r + \frac{1}{|r - R|} \right) \Psi(r) = -E \Psi(r)$$

3D partial differential equation!

$$Y_0^0 = 1$$



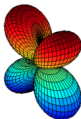
$$Y_1^0 = \cos\theta$$



$$Y_2^0 = 3\cos^2\theta - 1$$



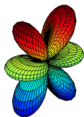
$${}^{\pm}Y_2^1 = \cos\theta \sin\theta \sin\phi$$



$$Y_3^0 = 5\cos^3\theta - 3\cos\theta$$



$${}^{\pm}Y_3^1 = (5\cos^2\theta - 1)\sin\theta \cos\phi$$



The hydrogen molecule and the Hartree-Fock method

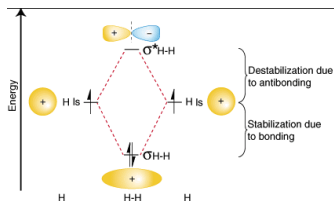
The Hamiltonian reads

$$H = -\frac{1}{2}\Delta_{r_1} - \frac{1}{|r_1 - R_1|} - \frac{1}{2}\Delta_{r_2} - \frac{1}{|r_2 - R_2|} \\ - \frac{1}{|r_2 - R_1|} - \frac{1}{|r_1 - R_2|} + \frac{1}{|r_1 - r_2|}$$

we want to solve

$$H\Psi(r_1, r_2) = E\Psi(r_1, r_2)$$

⇒ Not possible!



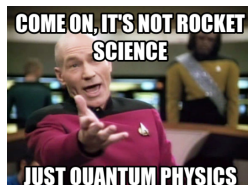
Second quantization I:

Fock space, creation and annihilation operators, CAR algebra

High-dimensional PDE's are hard!

- Can we think of matrices instead?
- Can we manipulate these matrices fast?

Second Quantization



Second quantization II:

Excitation operators, CCR algebra, nilpotent Lie algebra

The second quantization shows its teeth



Second quantization III:

Slater-Condon rules, Wick's theorem, Hartree-Fock revisited

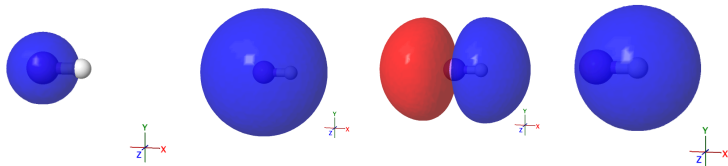
Let's bring out the big guns...



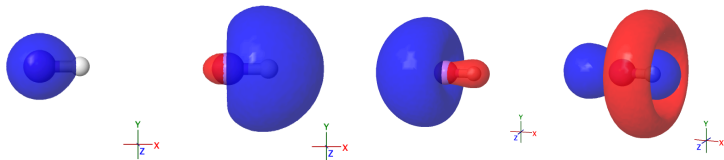
(Wick'ed guys)

The linear combination of atomic orbitals ansatz

Atomic orbitals



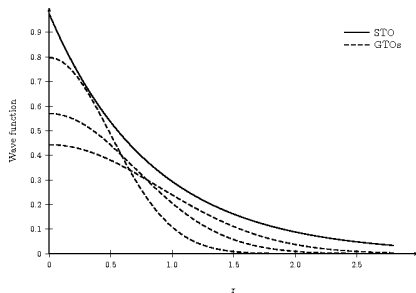
Molecular orbitals



Gaussian-type atomic orbitals and high-dimensional integration

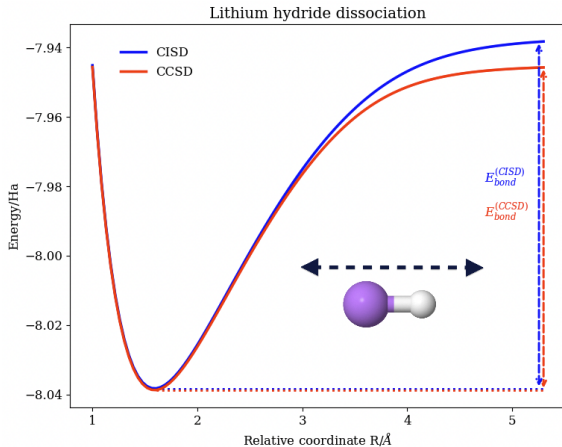
We will face the following integral evaluation:

$$v_{p,q,r,s} = \int_{X \times X} \frac{\xi_p^*(x_1)\xi_q(x_1)\xi_r^*(x_2)\xi_s(x_2)}{|r_1 - r_2|} d\lambda(x)d\lambda(x),$$



Size consistency and the coupled cluster ansatz

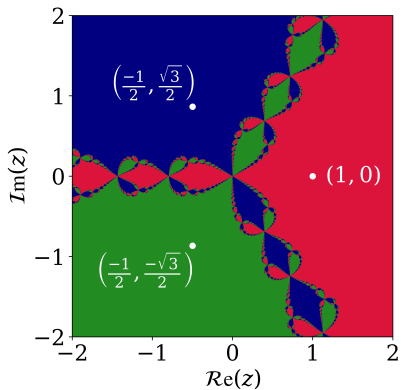
What do we really care about?



Truncating the coupled cluster ansatz and its working equations

$$\begin{aligned}
 E(t) &= \langle \Psi_0, H \Psi_0 \rangle + \sum_{IA} f_{IA} t_I^A + \frac{1}{4} \sum_{IJA B} \langle IJ \| AB \rangle t_{IJ}^{AB} + \frac{1}{2} \sum_{IJA B} \langle IJ \| AB \rangle t_I^A t_J^B, \\
 J(t)^A &= f_{IA} + \sum_C f_{AC} t_C^C - \sum_K f_{KI} t_K^A + \sum_{KC} \langle KA \| CI \rangle t_C^K + \sum_{KC} f_{KC} t_I^A t_C^K \\
 &\quad + \frac{1}{2} \sum_{KCD} \langle KA \| CD \rangle t_{KI}^{CD} - \frac{1}{2} \sum_{KLC} \langle KL \| CI \rangle t_{KL}^A - \sum_{KC} f_{KC} t_I^C t_K^A - \sum_{KLC} \langle KL \| CI \rangle t_C^K t_L^A \\
 &\quad + \sum_{KCD} \langle KA \| CD \rangle t_K^C t_I^D - \sum_{KLCD} \langle KL \| CD \rangle t_K^C t_I^D t_L^A + \sum_{KLCD} \langle KL \| CD \rangle t_C^K t_L^A t_I^D \\
 &\quad - \frac{1}{2} \sum_{KLCD} \langle KL \| CD \rangle t_{KI}^C t_L^A - \frac{1}{2} \sum_{KLCD} \langle KL \| CD \rangle t_{KL}^C t_I^D \\
 J(t)^{AB} &= \langle IJ \| AB \rangle + \sum_C \langle JBC \rangle t_I^A t_J^C - \langle JAC \rangle t_I^B t_J^C - \sum_K \langle JKC \rangle t_I^A t_K^B - \langle JKI \rangle t_J^B t_K^A \\
 &\quad + \frac{1}{2} \sum_{KL} \langle KL \| IJ \rangle t_{KL}^{AB} + \frac{1}{2} \sum_{CD} \langle AB \| CD \rangle t_{IJ}^{CD} + P(IJ)P(AB) \sum_{KC} \langle KB \| CJ \rangle t_{IK}^{AC} \\
 &\quad + P(IJ) \sum_C \langle AB \| CJ \rangle t_I^C - P(AB) \sum_K \langle KB \| IJ \rangle t_K^A \\
 &\quad + \frac{1}{2} P(IJ)P(AB) \sum_{KLCD} \langle KL \| CD \rangle t_{IK}^A t_{LJ}^{DB} + \frac{1}{4} \sum_{KLCD} \langle KL \| CD \rangle t_{IJ}^C t_{KL}^{AB} \\
 &\quad + \frac{1}{2} P(AB) \sum_{KLCD} \langle KL \| CD \rangle t_{IJ}^A t_{KL}^{BD} - \frac{1}{2} P(IJ) \sum_{KLCD} \langle KL \| CD \rangle t_{IK}^A t_{JL}^{CD} \\
 &\quad + \frac{1}{2} P(AB) \sum_{KL} \langle KL \| IJ \rangle t_K^A t_L^B + \frac{1}{2} P(IJ) \sum_{CD} \langle AB \| CD \rangle t_I^C t_J^D \\
 &\quad - P(IJ)P(AB) \sum_{KC} \langle KB \| IC \rangle t_K^A t_I^C + P(AB) \sum_{KC} f_{KC} t_K^A t_I^B t_C^C \\
 &\quad + P(IJ) \sum_{KC} f_{KC} t_I^C t_K^A t_J^B - P(IJ) \sum_{KLC} \langle KL \| CI \rangle t_K^C t_L^A t_I^B \\
 &\quad + P(AB) \sum_{KCD} \langle KA \| CD \rangle t_K^A t_I^B t_J^C + P(IJ)P(AB) \sum_{KCD} \langle AK \| DC \rangle t_I^D t_J^B t_C^K \\
 &\quad + P(IJ)P(AB) \sum_{KLC} \langle KL \| IC \rangle t_L^A t_J^B t_C^K + \frac{1}{2} P(IJ) \sum_{KLC} \langle KL \| CJ \rangle t_I^C t_K^A t_L^B \\
 &\quad - \frac{1}{2} P(AB) \sum_{KCD} \langle KB \| CD \rangle t_K^A t_I^C t_J^D + \frac{1}{2} P(IJ)P(AB) \sum_{KLC} \langle KB \| CD \rangle t_I^C t_K^A t_L^D \\
 &\quad + \frac{1}{2} P(IJ)P(AB) \sum_{KLC} \langle KL \| CJ \rangle t_I^C t_K^A t_L^B - P(IJ) \sum_{KLCD} \langle KL \| CD \rangle t_C^K t_I^D t_L^A t_J^B \\
 &\quad - P(AB) \sum_{KLCD} \langle KL \| CD \rangle t_C^K t_I^D t_J^B t_L^A - \frac{1}{4} P(IJ) \sum_{KLCD} \langle KL \| CD \rangle t_I^C t_J^D t_K^A t_L^B \\
 &\quad + \frac{1}{4} P(AB) \sum_{KLCD} \langle KL \| CD \rangle t_K^A t_L^B t_I^C t_J^D + P(IJ)P(AB) \sum_{KLCD} \langle KL \| CD \rangle t_I^C t_L^B t_K^A t_J^D \\
 &\quad + \frac{1}{4} P(IJ)P(AB) \sum_{KLCD} \langle KL \| CD \rangle t_I^C t_K^A t_J^D t_L^B
 \end{aligned}$$

Finding one root to the coupled cluster equations



Finding all roots to the coupled cluster equations

Seeking:

$$f(\mathbf{t}) = \begin{bmatrix} f_1(\mathbf{t}) \\ \vdots \\ f_m(\mathbf{t}) \end{bmatrix} = \begin{bmatrix} f_1(t_1, \dots, t_m) \\ \vdots \\ f_m(t_1, \dots, t_m) \end{bmatrix} = 0.$$

Find a start system g , and then continuously deform g into f

Davidenko differential equation:

$$\frac{\partial}{\partial \mathbf{t}} H(\mathbf{t}, \lambda) \left(\frac{d}{d\lambda} \mathbf{t}(\lambda) \right) + \frac{\partial}{\partial \lambda} H(\mathbf{t}, \lambda) = 0, \quad \mathbf{t}(1) = \mathbf{s}_0,$$

Finding all roots to the coupled cluster equations

Seeking:

$$f(t) = \begin{bmatrix} f_1(t) \\ \vdots \\ f_m(t) \end{bmatrix} = \begin{bmatrix} f_1(t_1, \dots, t_m) \\ \vdots \\ f_m(t_1, \dots, t_m) \end{bmatrix} = 0.$$

